

Math 678.
Lecture 9.

IVP:
$$\begin{cases} u_t - \Delta u = 0 & \text{in } \mathbb{R}^n \times (0, \infty) \\ u = g & \text{on } \mathbb{R}^n \times \{t=0\} \end{cases}$$

$$u(x, t) = \int_{\mathbb{R}^n} \Phi(x-y, t) g(y) dy \quad (*)$$

Thm: $g \in C(\mathbb{R}^n) \cap L^\infty(\mathbb{R}^n) \Rightarrow (*)$ satisfies:

- 1) $u \in C^\infty(\mathbb{R}^n \times (0, \infty))$
- 2) $u_t = \Delta u, t > 0$
- 3) $\lim_{\substack{x \rightarrow x^0, t \rightarrow 0 \\ t > 0}} u(x, t) = g(x^0), x^0 \in \mathbb{R}^n$

Nonhomogeneous heat eqn:

$$\begin{cases} u_t - \Delta u = f, & \text{in } \mathbb{R}^n \times (0, \infty) \\ u = g & \text{on } \mathbb{R}^n \times \{t=0\} \end{cases}$$

Duhamel's principle

1) In linear algebra: ~~let~~ let u_0, u_1, \dots in \mathbb{R}^n

satisfy $u_{j+1} = A_{j+1} u_j$

A_1, A_2, \dots - sequence of operators on \mathbb{R}^n .

$$u_k = A_k A_{k-1} \dots A_1 u_0$$

$$u_k = \underbrace{A_k A_{k-1} \dots A_{j+1}}_{S(k, j)} u_j$$

if $k > j$ $S(k, j) \quad S(1, 1) = I$

$$\Rightarrow u_k = S(k, j) u_j$$

Nonhomogeneous problem: $u_{j+1} = A_{j+1} u_j + v_{j+1}$

v_1, v_2, \dots - sequence of given vectors

$$u_1 = A_1 u_0 + v_1 = S(1, 0) u_0 + v_1$$

$$\begin{aligned} u_2 &= A_2 u_1 + v_2 = A_2 (A_1 u_0 + v_1) + v_2 = \\ &= A_2 A_1 u_0 + A_2 v_1 + v_2 \\ &= S(2, 0) u_0 + S(2, 1) v_1 + v_2 \end{aligned}$$

$$\begin{aligned} u_3 &= A_3 u_2 + v_3 = A_3 (A_2 A_1 u_0 + A_2 v_1 + v_2) + v_3 \\ &= A_3 A_2 A_1 u_0 + A_3 A_2 v_1 + A_3 v_2 + v_3 \\ &= S(3, 0) u_0 + S(3, 1) v_1 + S(3, 2) v_2 \\ &\quad + v_3 \end{aligned}$$

$$u_k = S(k, 0) u_0 + S(k, 1) v_1 + S(k, 2) v_2 + \dots \Rightarrow$$

$$u_k = S(k, 0) u_0 + \sum_{j=1}^k S(k, j) v_j$$

2) In ODE: $\frac{du}{dt} = A(t)u(t)$, $t > 0$ (1)

if $u(t) = S(t, t_0) u(t_0)$ solves (1) with $u(0) = g$.
and $A(t)$ - continuous map

Then solution to $\begin{cases} \frac{du}{dt} = A(t)u(t) + f(t), t > 0 \\ u(0) = g \end{cases}$

is given by $u(t) = S(t, 0) u(0) + \int_0^t S(t, s) f(s) ds$

Justification: $\frac{d}{dt} (S(t, t_0) u(t_0)) = A(t) \cdot S(t, t_0) u(t_0)$
 ~~$\frac{du}{dt} =$~~ $\frac{d}{dt} (S(t, 0) u(0)) = A(t) \cdot S(t, 0) u(0)$

$$\Rightarrow \frac{du}{dt} = A(t) S(t, 0) u(0) + \int_0^t A(t) S(t, s) f(s) ds + \underbrace{S(t, t)}_I \underbrace{f(t)}_{\int^t \frac{df}{ds} ds}$$

$$\Rightarrow \frac{du}{dt} = A(t) \cdot \left[S(t, 0) u(0) + \int_0^t S(t, s) f(s) ds \right] + f(t)$$

$$\frac{du}{dt} = A(t) u(t) + f(t)$$

3) In PDE case:

$$u = u(x, t; s) = \int_{\mathbb{R}^n} \Phi(x-y, t-s) f(y, s) dy$$

$$\text{solves } \begin{cases} u_t(\cdot; s) - \Delta u(\cdot; s) = 0 & , \text{ in } \mathbb{R}^n \times (s, \infty) \\ u(\cdot; s) = f(\cdot; s) & \text{ on } \mathbb{R}^n \times \{t=s\} \end{cases}$$

IVP with time starting at $t=s$

Consider the non-homogeneous problem:

$$\begin{cases} u_t - \Delta u = f, & \text{in } \mathbb{R}^n \times (0, \infty) \\ u = 0 & \text{on } \mathbb{R}^n \times \{t=0\} \end{cases} \quad (**)$$

By Duhamel's principle we obtain $u(x, t)$ -
Solution to $(**)$ - by integrating $u(x, t; s)$:

$$u(x, t) = \int_0^t u(x, t; s) ds, \quad x \in \mathbb{R}^n, t > 0$$

$$= \int_0^t \int_{\mathbb{R}^n} \Phi(x-y, t-s) f(y, s) dy ds =$$

$$= \int_0^t \frac{1}{(4\pi(t-s))^{n/2}} \cdot \int_{\mathbb{R}^n} e^{-\frac{|x-y|^2}{4(t-s)}} f(y, s) dy ds$$

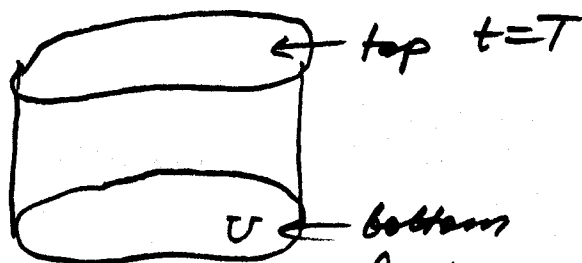
Consequence: $\begin{cases} u_t - \Delta u = f & \text{in } \mathbb{R}^n \times (0, \infty) \\ u = g & \text{on } \mathbb{R}^n \times \{t=0\} \end{cases}$

$$\Rightarrow u(x, t) = \int_{\mathbb{R}^n} \Phi(x-y, t) g(y) dy + \int_0^t \int_{\mathbb{R}^n} \Phi(x-y, t-s) f(y, s) dy ds$$

Mean Value formula.

$U_T := U \times (0, T]$ - parabolic cylinder

$\Gamma_T := \overline{U_T} - U_T$ - parabolic boundary = bottom + sides



U_T - parabolic interior - includes the top.

$\Phi(x-y)$ - fund. sol. of Laplace is constant on any sphere $\partial B(x,r)$
level set of Φ

$$E(x,t;r) := \left\{ (y,s) \in \mathbb{R}^{n+1} \mid s \leq t, \Phi(x-y, t-s) \geq \frac{1}{r^n} \right\}$$

"heat ball"

∂E is the level set for Φ ↑
fund. sol. of heat eqn



Thm: (MVT for heat eqn)

$u \in C_1^2(U_T)$ - the soln to heat eqn

$$\text{Then } u(x,t) = \frac{1}{4r^n} \iint_{E(x,t;r)} u(y,s) \frac{|x-y|^2}{(t-s)^2} dy ds$$

$\forall E(x,t;r) \subset U_T$

Proof: Change coordinates to $x=0, t=0$.

$$\Rightarrow E(r) = E(0,0;r)$$

Assume u is smooth (otherwise use a mollifier)

$$\varphi(r) := \frac{1}{r^n} \int_{E(r)} u(y,s) \frac{|y|^2}{s^2} dy ds = \iint_{E(r)} u(ry, r^2s) \frac{|y|^2}{s^2} dy ds$$

$$\varphi'(r) = \iint_{E(r)} \left(\sum_{i=1}^n u_{y_i} \cdot y_i \cdot \frac{|y|^2}{s^2} + 2r \cdot u_s \cdot \frac{|y|^2}{s^2} \right) dy ds$$

$$= \frac{1}{r^{n+1}} \iint_{E(r)} \left(\underbrace{\sum u_{y_i} \cdot y_i \cdot \frac{|y|^2}{s^2}}_I + \underbrace{2u_s \frac{|y|^2}{s^2}}_J \right) dy ds$$

$$\Psi := \frac{-n}{2} \log(-4\pi s) + \frac{|y|^2}{4s} + n \log r$$

(1) $\Psi = 0$ on $\partial E(r)$ since ∂E has $\Phi(-y, -s) = \frac{1}{r^n}$
 on $\partial E(r)$: $\Phi(y, -s) =$

$$= \frac{1}{(-4\pi s)^{n/2}} e^{+|y|^2/4s} = \frac{1}{r^n}$$

$$\Phi(x, t) = \begin{cases} \frac{1}{(4\pi t)^{n/2}} e^{-\frac{|x|^2}{4t}}, & t > 0 \\ 0, & t < 0 \end{cases}$$

$$e^{+|y|^2/4s} = \frac{(-4\pi s)^{n/2}}{r^n}$$

$$\frac{|y|^2}{4s} = \frac{n}{2} \log(-4\pi s) - n \log r \text{ on } \partial E(r)$$

$$\Rightarrow \Psi = -\frac{n}{2} \log(-4\pi s) + \frac{|y|^2}{4s} + n \log r = 0 \text{ on } \partial E(r)$$

(2) $\frac{\partial \Psi}{\partial y_i} = \frac{y_i}{2s} \Rightarrow$

Now $I = \frac{1}{r^{n+1}} \iint_{E(r)} \sum u_{y_i} \cdot y_i \cdot \frac{|y|^2}{4s} dy ds =$

$$= \frac{1}{r^{n+1}} \iint 4u_s \sum u_{y_i} \cdot \Psi y_i dy ds$$